

Low scale strings and dijets at the LHC

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- **D-branes and all that**
- **Regge vs. KK or winding modes**
- **Dijet signals at the LHC**

**L. Anchordoqui, HG, D. Lüster, S. Nawata, S. Stieberger, T. Taylor,
PRL 101, 241803 (2008)**

TeV scale strings

- Large extra spatial dimensions and D-brane constructs allow
 - low string scale compatible with weak 4-D gravity – in toroidal compactification

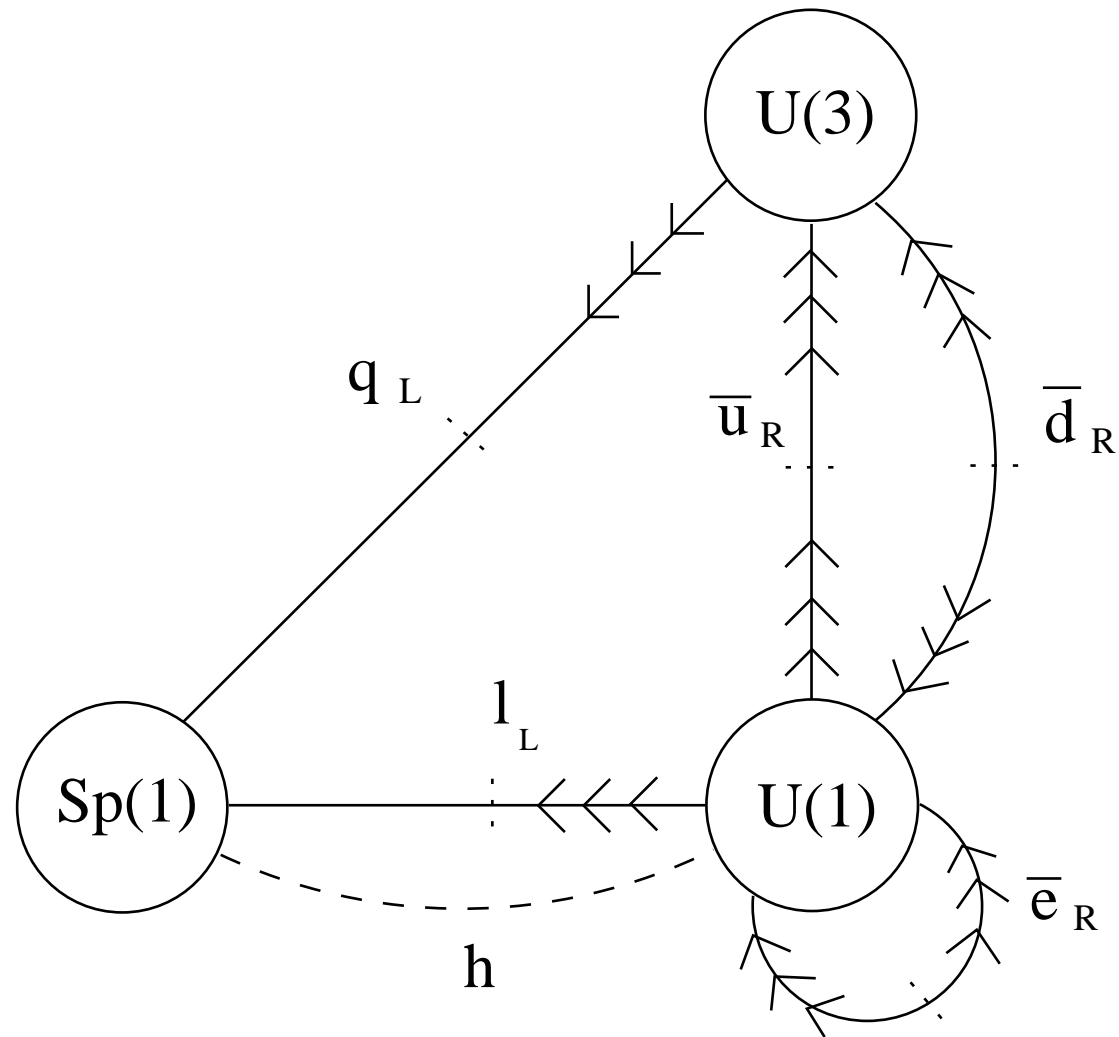
$$M_{\text{Pl}}^2 \sim M_s^2 (M_s R)^n$$

- Regge recurrences in TeV region
- Open strings can terminate on stack of N identical D branes
 - $U(N)$ gauge group for each stack if images under orientifolding are different from themselves

TeV scale strings (cont'd)

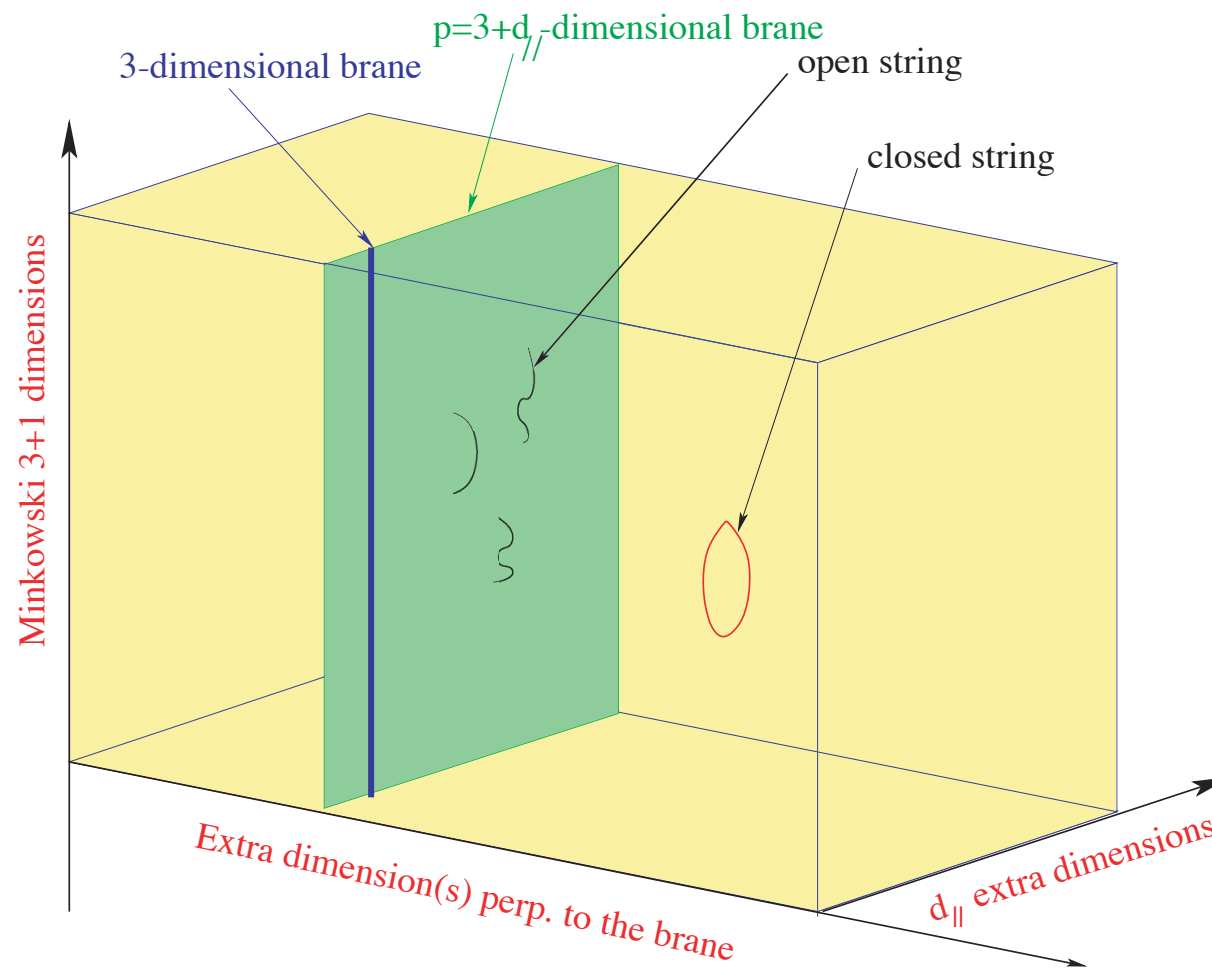
- $Sp(N)$ or $SO(N)$ gauge group for each stack if they self-image under orientifolding
- Matter fields are in bifundamental representations (N_a, N_b) or (N_a, \bar{N}_b) located at intersections of D-branes
- Regge vs. KK/winding modes

Quivers



Berenstein, Pinansky, [hep-th/0610104]; Antoniadis, Kiritsis, Tomaras, [hep-ph/0004214]

Parallel and perpendicular extra dimensions



Antoniadis [[hep-th/0710.4267](https://arxiv.org/abs/hep-th/0710.4267)]

Gauge fields

- $U(3)$: 8 $SU(3)$ gluons, additional $U(1)$ (C_μ) coupled to baryon number with strength $g_3/\sqrt{6}$
- $U(2)$: 3 $SU(2)$ W' 's, additional $U(1)$
- **Minimal:** $Sp(1)$: 3 W' 's, no additional $U(1)$
- $U(1)$: another $U(1)$ (B_μ)

Y_μ (hypercharge) = linear combination of C_μ and other $U(1)$'s

Charge assignments, anomalies

- With Q_B assignments $(0, -1, 1, 1, 2)$ for $(q_L, u_R, d_R, \ell_L, e_R)$, the hypercharge

$$Q_Y = \frac{1}{6} Q_{U(3)} - \frac{1}{2} Q_{U(1)}$$

is free of gauge and mixed anomalies.

- The orthogonal charges $Q_{Y'}$ suffer from anomalies
- Cancel via Green-Schwartz mechanism - **coupling** $\eta F \wedge F$ to RR closed string two-form field

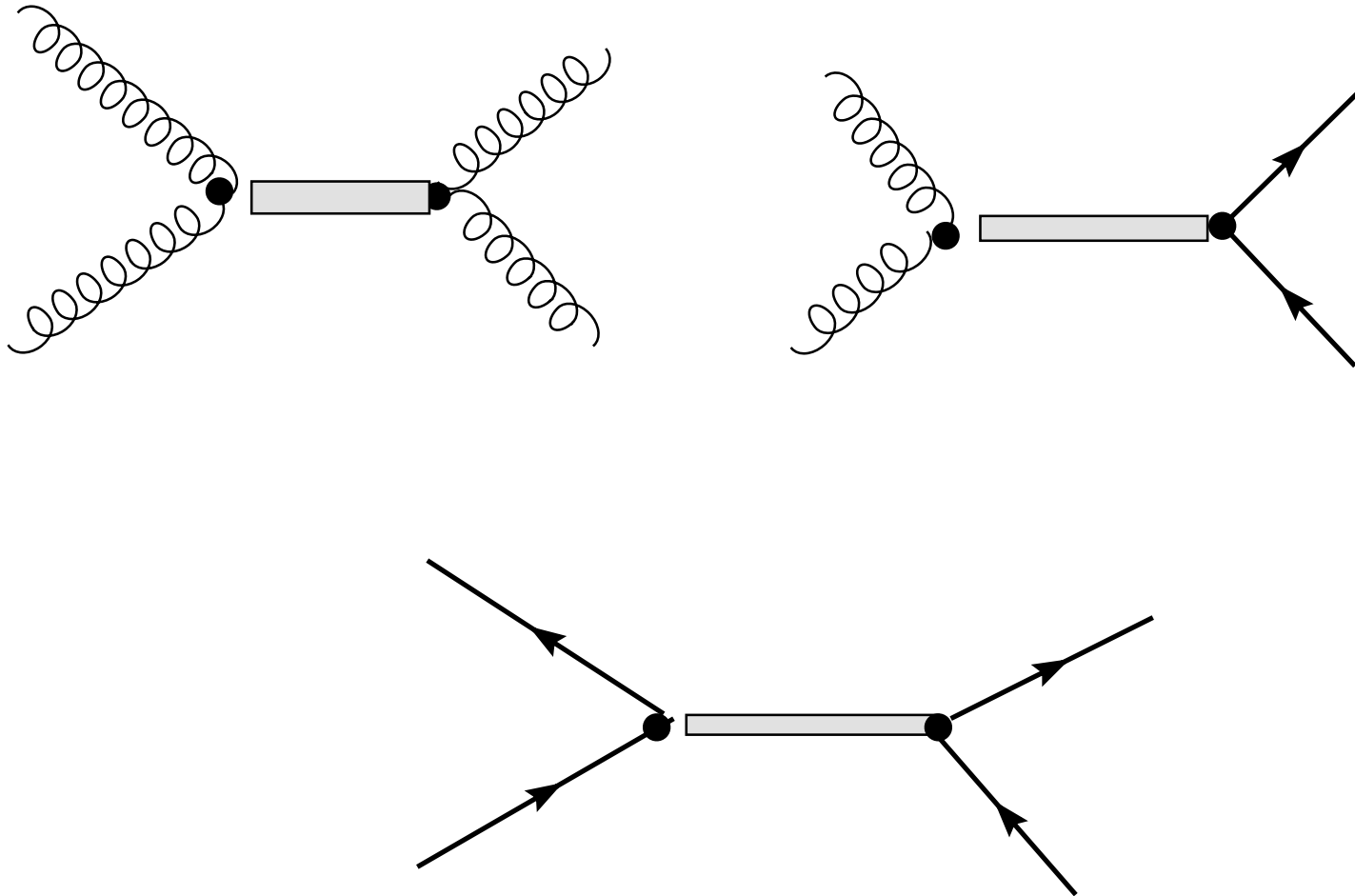
Gauging baryon number

- If left unbroken, long range force coupled to baryon number – violation of everything
- If broken via a Higgs mechanism, break global baryon number at TeV scale → **fast proton decay**
- Mixing $B \wedge F$ with closed string two-form $B_{\mu\nu}$ → mass for Y' , **global baryon number preserved.**
Akin to Stückelberg mechanism
- Allows TeV BH production in νp collisions w/out violation of baryon number via BH

Model independence on the color brane

- Consider scattering on the $U(3)$ brane, involving quarks and gluons
- Gluons live only on a single stack of color branes, don't see other branes
- Momentum in all the $p - 3$ extra parallel dimensions conserved, so **4-point functions** $gg \rightarrow gg$, $q\bar{q} \rightarrow gg$, $qg \rightarrow qg$ do not excite KK or winding modes – **only Regge recurrences independent of compactification.**
- Not so for $qq \rightarrow qq$ – these live on intersecting branes, and parallel momentum not conserved – **KK/winding modes involved.**

QCD processes



Amplitudes

The basic string amplitude for one ordering (out of six) of 4 $U(3)$ gauge bosons is

$$A(1^-, 2^-, 3^+, 4^+) = 4g^2 \text{Tr} (T^{a_1} T^{a_2} T^{a_3} T^{a_4}) \cdot \frac{\langle 12 \rangle^4}{\langle 12 \rangle \langle 23 \rangle \langle 34 \rangle \langle 41 \rangle} V(k_1, k_2, k_3, k_4)$$

where

$$V(k_1, k_2, k_3, k_4) = \frac{\Gamma(1-s) \Gamma(1-u)}{\Gamma(1+t)}$$

$$\langle ij \rangle = \bar{u}_L(k_i) u_R(k_j)$$

S. Stieberger and T. R. Taylor, hep-th/0609175

Notes

- Only one basic amplitude (MHV, or Maximum Helicity Violating) Parke and Taylor, PRL 1986
- Origin is relation to single 4 scalar amplitude in $N = 2$ supersymmetry
- Obtained as scattering of open strings: insertion of 4 vector vertex functions on the boundary of the disk, which is projection of upper half plane of world sheet complex coordinate
- Similarly for $ggq\bar{q}$.

Resonant cross sections

- Go to resonant poles at $s = M_s^2$:
 - bosonic** g^* , C^* , $J = 0, 2$;
 - fermionic** q^* , $J = 1/2, 3/2$
- Soften to Breit Wigner form, obtaining vertex functions and widths by factorizing amplitudes at the pole. **Example:**

$$\begin{aligned} \Gamma^{J=2}(g^* \rightarrow gg) &= (1/6) (g^2/4\pi) M \\ \Gamma^{J=2}(g^* \rightarrow gC) &= (2/15) (g^2/4\pi) M \\ \Gamma^{J=2}(g^* \rightarrow q\bar{q}) &= (3/20) (g^2/4\pi) M \end{aligned}$$

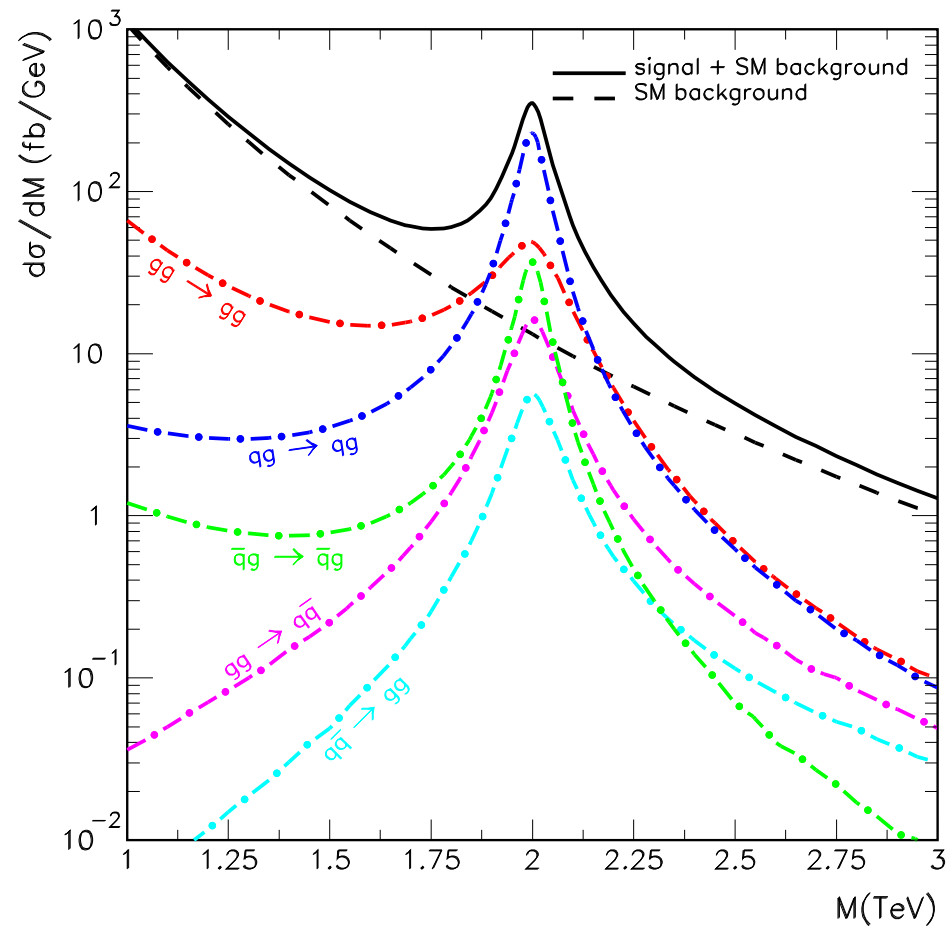
L. A. Anchordoqui, HG and T. R. Taylor, Phys. Lett. B 668, 373 (2008)

Example

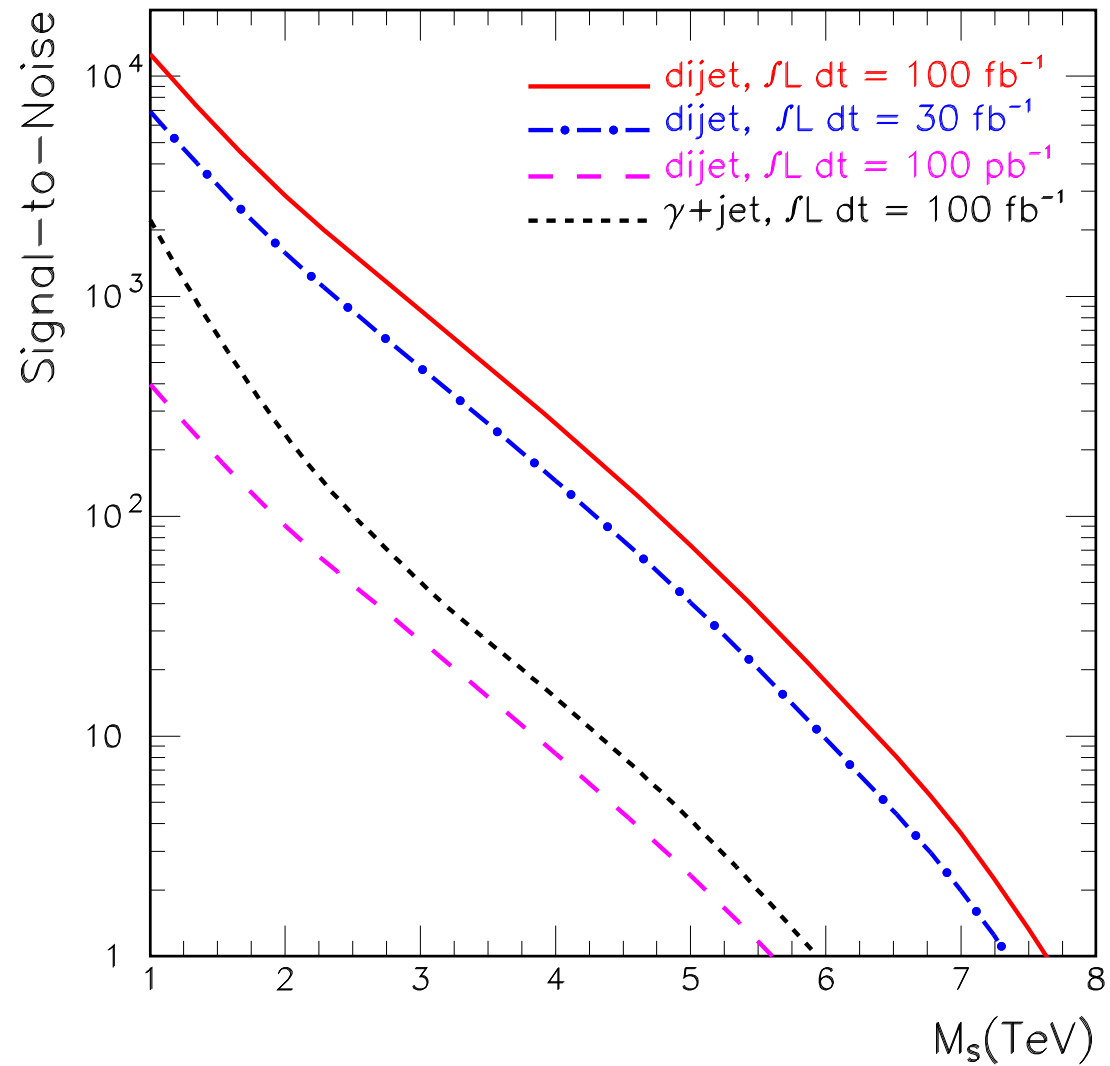
$$\begin{aligned}
 |\mathcal{M}(gg \rightarrow gg)|^2 &= \frac{19}{12} \frac{g^4}{M_s^4} \left\{ \right. \\
 &W_{g^*}^{gg \rightarrow gg} \left[\frac{M_s^8}{(s - M_s^2)^2 + (\Gamma_{g^*}^{J=0} M_s)^2} \right. \\
 &\quad \left. + \frac{t^4 + u^4}{(s - M_s^2)^2 + (\Gamma_{g^*}^{J=2} M_s)^2} \right] \\
 &+ W_{C^*}^{gg \rightarrow gg} \left[\frac{M_s^8}{(s - M_s^2)^2 + (\Gamma_{C^*}^{J=0} M_s)^2} \right. \\
 &\quad \left. + \frac{t^4 + u^4}{(s - M_s^2)^2 + (\Gamma_{C^*}^{J=2} M_s)^2} \right] \left. \right\}
 \end{aligned}$$

Cuts, $d\sigma/dM$

- Set (**parton**)jet rapidity cuts $|y_1|, |y_2| < 1$
transverse momenta $p_{1,2}^T > 50 \text{ GeV}$



Discovery reach at LHC



Results

- Within a few years of data collection, **string scales up to 6.8 TeV** open to discovery at 5σ
- Very early in the initial run (100 pb^{-1}) **resonant mass as large 4 TeV** observable at 10σ

previous discussion of dijet TeV-scale string resonances at LHC, see P. Burikham, T. Figy, T. Han, e-Print: [hep-ph/0411094](https://arxiv.org/abs/hep-ph/0411094)

Angular distributions: the R parameter

B. Abbott *et al* [DØ Collaboration] PRL 82, 2457 (1999)

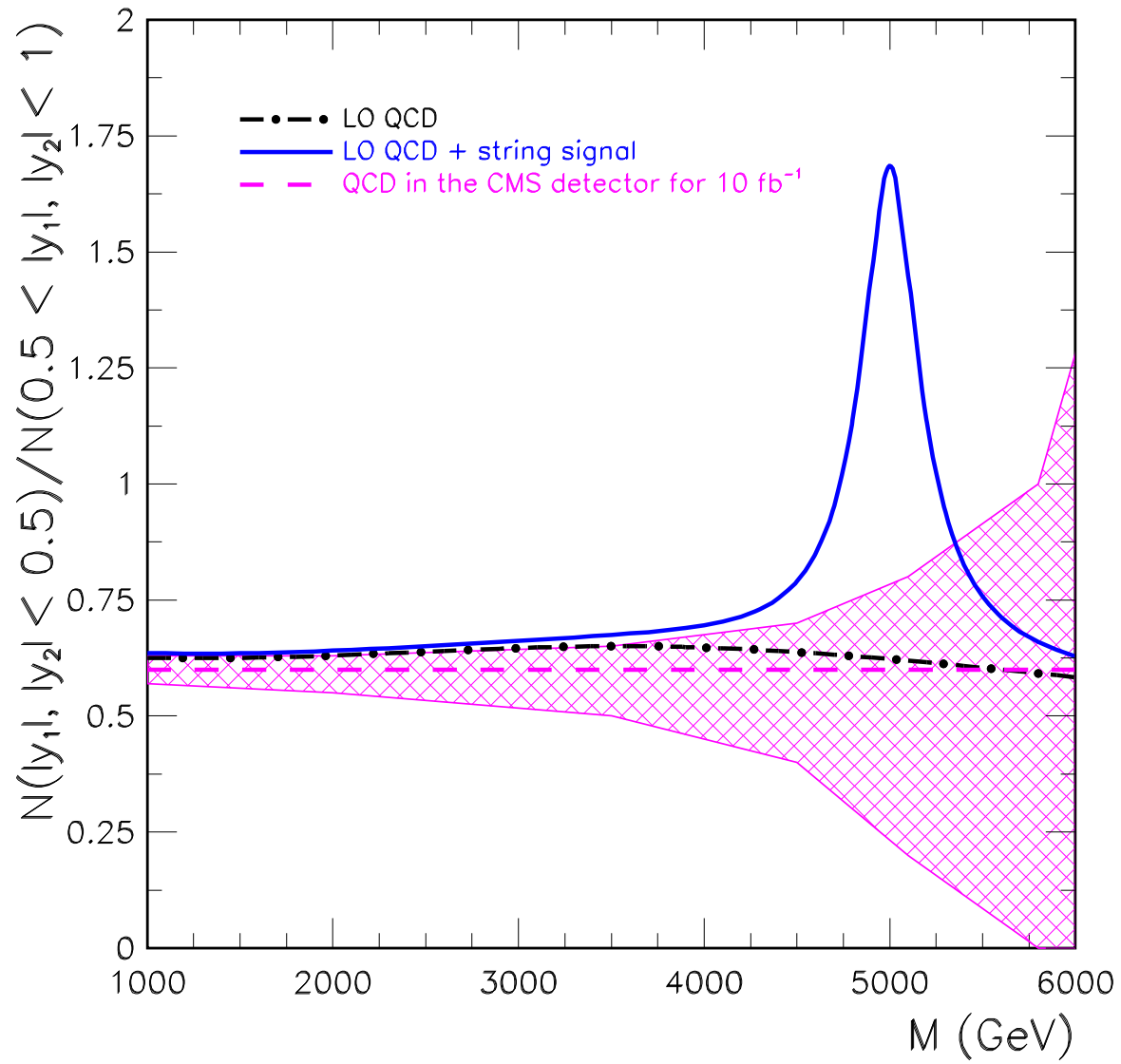
- QCD parton-parton cross sections dominated by t -channel exchanges \rightarrow **dijets at large rapidity**
- Non-standard contact interactions or excitations of resonances \rightarrow more isotropic distribution \rightarrow **more uniform distribution in rapidity.**
- **Define**

$$R = \frac{N_{\text{ev}}(|y_1|, |y_2| < 0.5)}{N_{\text{ev}}(0.5 < |y_1|, |y_2| < 1.0)}$$

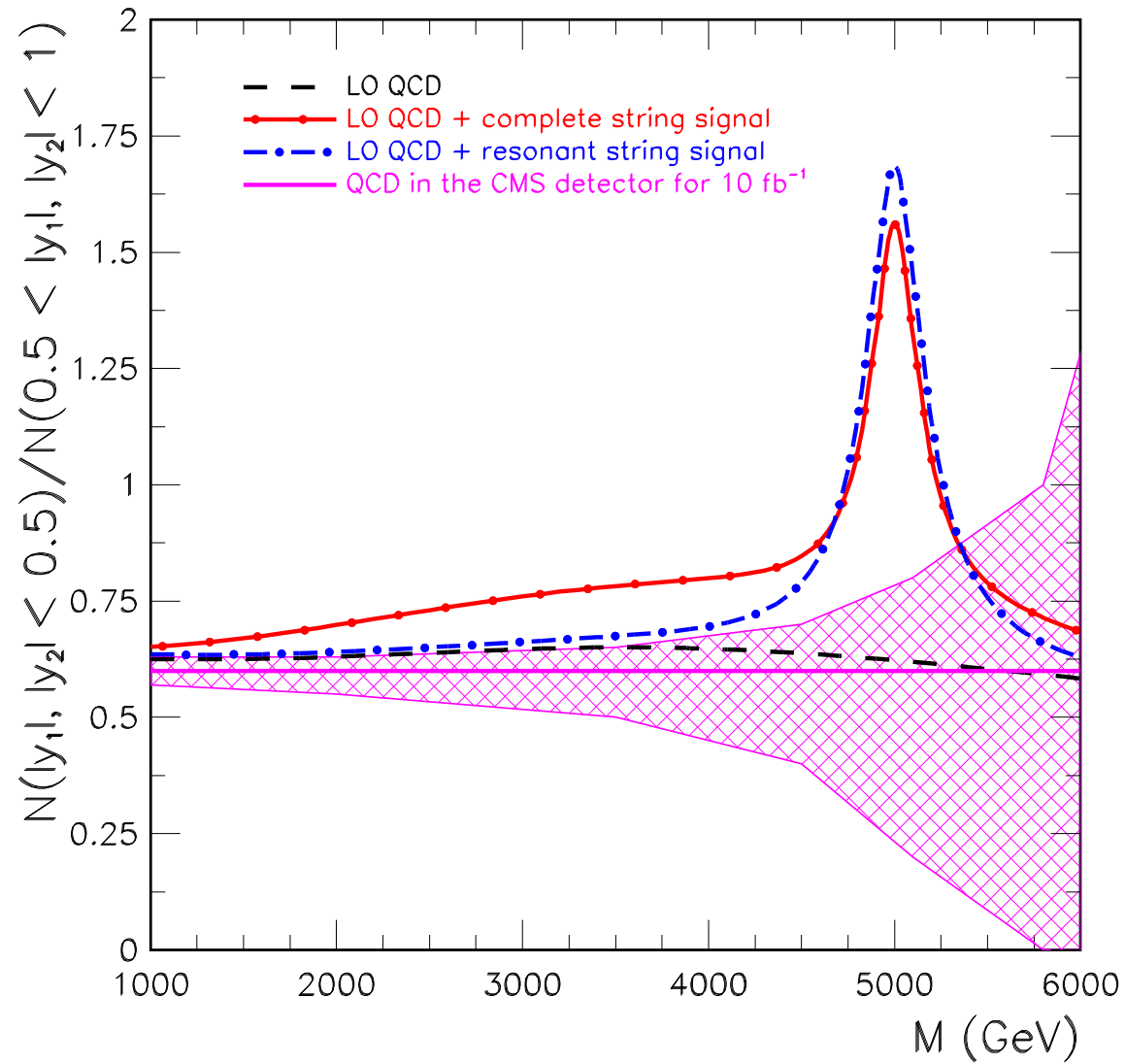
- **For QCD $R \simeq 0.6$ independent of M**

For previous discussion of R parameter in string theory dijets, see P.Meade and L.Randall, arXiv:0708.3017

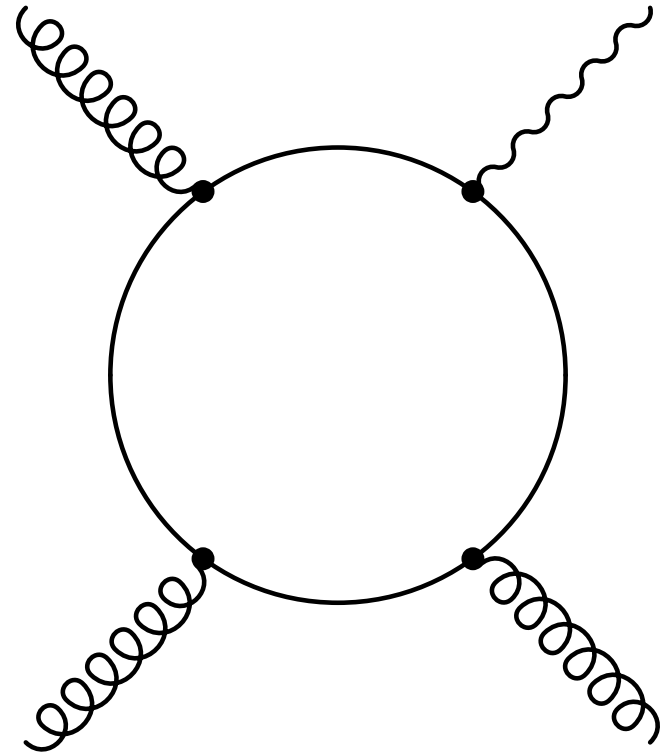
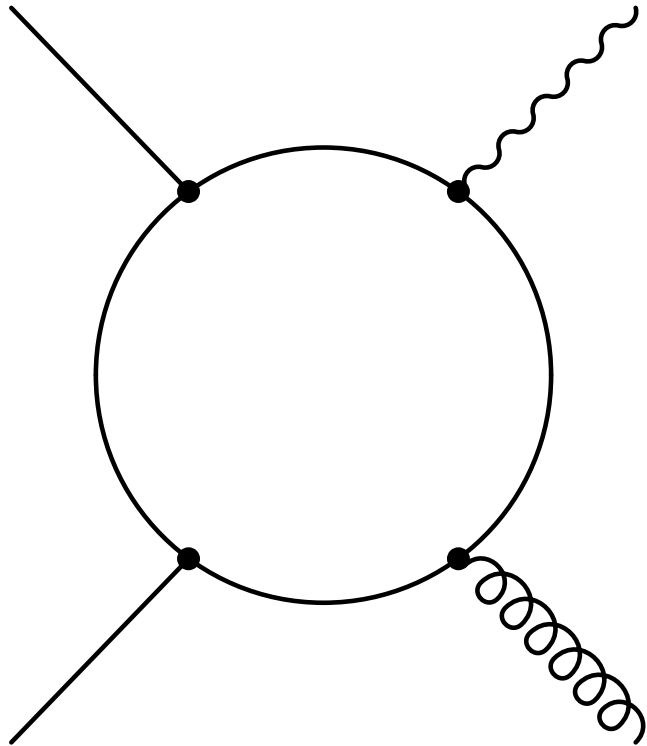
Regge in the R plot



qq in the R plot



Photon processes



$$gg \rightarrow \gamma g$$

L. Anchordoqui et al, PRL 100, 171603 (2008); PR D78, 016005 (2008)

- Does not exist at tree level in field theory
- But does exist at tree (i.e. disk) level in string theory: **U(1) field on color stack has photon component**
- Involves only gauge bosons on a single stack, so is **independent of fermion embeddings on the intersecting stacks**
- Need to impose isolation procedure to greatly reduce π^0 background

Conclusions

- We have calculated the discovery reach at LHC for TeV-scale open string theory via the mass spectrum and angular distribution of dijets
- The key provided by **quiver assignments** – **gauge supermultiplets** on single stack of D-branes, **matter supermultiplets** bifundamental assignments at intersecting stacks
- In this picture, for all the $2 \rightarrow 2$ partonic subprocesses, **except for $qq \rightarrow qq$, $q\bar{q} \rightarrow q\bar{q}$** , the string amplitude **is completely determined** –
 - **no arbitrary Chan-Paton factors**
 - **no dependence on compactification scheme**
 - **no dependence on fermion embeddings**

Conclusions (concluded)

- Even with the 4-fermion processes **omitted from the signal but included in the QCD background**, the resonant signal for the first Regge recurrence can attain **discovery at 5σ** for strings scales as high as **6.8 TeV**
- Discovery is greatly enhanced through observation of angular distributions as manifest in the R parameter
- Observation of **$\gamma + \text{jet}$** can provide independent corroborating evidence if the string scale is smaller than 5 TeV