

EFFECTIVE FIELD THEORY $T_{\mu\nu}$
AND THE
COSMOLOGICAL CONSTANT PROBLEM

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COSMOLOGICAL CONSTANT PROBLEMS

From QFT : Zero Point Energies contribution to CC

From QFT : Condensates contribution to CC

Why ρ_Λ has the tiny value we measure?

Why ρ_Λ and ρ^{mat} at present time coincide?

.... Modified gravity, Voids,

within OUR MODEL we address the (QFT) Problems 1 & 2

Conclusion: ZPE & Condensates do not contribute to CC.

Where does this come from?

From the Effective Nature of Field Theories

VACUUM ENERGY & CC

QUESTION : How do we usually manage with the
Zero Point Energies contribution to CC?

ANSWER :

1. Lorentz Invariance $\implies T_{\mu\nu}^{vac} = \rho^{vac} g_{\mu\nu} \quad (p^{vac} = -\rho^{vac})$

Now take the example of : $\mathcal{L} = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{1}{2} m^2 \phi^2$

2. Vacuum Energy Density (\Leftarrow Zero Point Energy) :

$$\rho^{vac} = \frac{1}{V} \sum_{\vec{k}} \frac{1}{2} \sqrt{\vec{k}^2 + m^2} = \frac{1}{2} \int \frac{d^3 \vec{k}}{(2\pi)^3} \sqrt{\vec{k}^2 + m^2} \simeq \frac{M_P^4}{16\pi^2}$$

3. Continuity Equation : $\dot{\rho} + 3 \left(\frac{\dot{a}}{a} \right) (\rho + p) = 0$

$$\rho^{vac} = Const. \simeq \frac{M_P^4}{16\pi^2} \geq 120 \text{ ord. of magn. problem ...}$$

..... BUT

EINSTEIN EQS. : ENERGY DENSITY & PRESSURE

... Let's start again from the beginning ...

$$G_{\mu\nu} - \lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$$

Which $T_{\mu\nu}$ in the Ein. Eq. from our $\mathcal{L} = \frac{1}{2}\partial_\mu\phi\partial^\mu\phi - \frac{1}{2}m^2\phi^2$?

Quantum Statistical Average (non diag. terms vanish):

$$T_{00} = \rho = \langle\langle \hat{T}_{00} \rangle\rangle = \frac{1}{V} \sum_{\vec{k}} \sum_n \langle n | \varrho | n \rangle n_{\vec{k}} \omega_{\vec{k}} + \frac{1}{V} \sum_{\vec{k}} \frac{1}{2} \omega_{\vec{k}}$$

$$T_{ii} = p = \langle\langle \hat{T}_{ii} \rangle\rangle = \frac{1}{V} \sum_{\vec{k}} \sum_n \langle n | \varrho | n \rangle n_{\vec{k}} \frac{(k^i)^2}{\omega_{\vec{k}}} + \frac{1}{V} \sum_{\vec{k}} \frac{(k^i)^2}{2\omega_{\vec{k}}}$$

$|n\rangle$: generic element Fock space basis;

ϱ : density operator ; $n_{\vec{k}} = \langle n | a_{\vec{k}}^\dagger a_{\vec{k}} | n \rangle$; $\omega_{\vec{k}} = \sqrt{\vec{k}^2 + m^2}$.

By performing the sum over n :

$$T_{00} = \rho = \int \frac{d^3 \vec{k}}{(2\pi)^3} \omega_{\vec{k}} n_{BE}(\vec{k}^2) + \frac{1}{2} \int \frac{d^3 \vec{k}}{(2\pi)^3} \omega_{\vec{k}} \equiv \rho^{mat} + \rho^{vac}$$

$$T_{ii} = p = \frac{1}{3} \int \frac{d^3 \vec{k}}{(2\pi)^3} \frac{\vec{k}^2}{\omega_{\vec{k}}} n_{BE}(\vec{k}^2) + \frac{1}{3} \int \frac{d^3 \vec{k}}{(2\pi)^3} \frac{\vec{k}^2}{2\omega_{\vec{k}}} \equiv p_{mat} + p^{vac}$$

$n_{BE}(\vec{k}^2)$ = Bose-Einstein distribution

In Blue : Gas of Rel. Part. (Matter/Radiation contribution)

In Red : The corresponding Vacuum contribution

Remark 1 : In the Einstein Eqs., ρ^{mat} & ρ^{vac} , as well as

p_{mat} & p^{vac} , enter on an equal footing.

Remark 2 : UV cutoff needed to compute ρ^{vac} & p^{vac}

On the contrary, due to the presence of $n_{BE}(\vec{k}^2)$, ρ^{mat} & p_{mat} are finite

Computing ρ^{vac} & p^{vac} with an UV cutoff $\Lambda (= M_P)$

for our $\mathcal{L} = \frac{1}{2}\partial_\mu\phi\partial^\mu\phi - \frac{1}{2}m^2\phi^2$ we have :

$$\rho^{vac} = \frac{1}{16\pi^2} \left[\Lambda(\Lambda^2 + m^2)^{\frac{3}{2}} - \frac{\Lambda m^2(\Lambda^2 + m^2)^{\frac{1}{2}}}{2} - \frac{m^4}{4} \ln \left(\frac{(\Lambda + (\Lambda^2 + m^2)^{\frac{1}{2}})^2}{m^2} \right) \right]$$
$$p^{vac} = \frac{1}{16\pi^2} \left[\frac{\Lambda^3(\Lambda^2 + m^2)^{\frac{1}{2}}}{3} - \frac{\Lambda m^2(\Lambda^2 + m^2)^{\frac{1}{2}}}{2} + \frac{m^4}{4} \ln \left(\frac{(\Lambda + (\Lambda^2 + m^2)^{\frac{1}{2}})^2}{m^2} \right) \right]$$

Note : For $\Lambda \gg m \implies$ RED dominant :

$$\rho^{vac} \simeq \frac{\Lambda^4}{16\pi^2} \quad ; \quad p^{vac} \simeq \frac{1}{3} \frac{\Lambda^4}{16\pi^2}$$

QUESTION : How do we deal with the Divergences ?

Formal Point of View : The Divergent Terms (which do not respect $T_{\mu\nu}^{vac} \propto g_{\mu\nu}$) have to be removed via **Renormalization**

(This would “alleviate” the 120 order of magnitude problem ...However, we still have the condensates...)

Criticism (De Witt) :

- **For Zero Point Energies : Physical Meaning of Divergences** rooted in the **harmonic oscillator** structure of a QFT.
- **Lost** if we cancel out these terms with a **formal procedure** such as **normal ordering**.
- Still, popular prescription for the automatic cancellation of these divergences : **Dimensional Regularization**

Deeper Physical Point of View \implies

Effective Field Theory Point of View

Lesson from Wilson RG & Effective Field Theories :

QFT = Effective Theory valid up to Λ ($\Lambda =$ “scale of new physics”)

Hierarchy of Field Theories up to M_P \leftarrow String theory (?)

According to this view, the cutoff Λ is physical \implies we do not discard any term in $T_{\mu\nu}$

OK. Let's take this point of view
and go back to the **Effective Field** $T_{\mu\nu}$
that we have computed before

Effective Field $T_{\mu\nu}$

$$T_{00} = \rho^{vac} = \frac{1}{16\pi^2} \left[\Lambda(\Lambda^2 + m^2)^{\frac{3}{2}} - \frac{\Lambda m^2(\Lambda^2 + m^2)^{\frac{1}{2}}}{2} - \frac{m^4}{4} \ln \left(\frac{(\Lambda + (\Lambda^2 + m^2)^{\frac{1}{2}})^2}{m^2} \right) \right]$$

$$T_{ii} = p^{vac} = \frac{1}{16\pi^2} \left[\frac{\Lambda^3(\Lambda^2 + m^2)^{\frac{1}{2}}}{3} - \frac{\Lambda m^2(\Lambda^2 + m^2)^{\frac{1}{2}}}{2} + \frac{m^4}{4} \ln \left(\frac{(\Lambda + (\Lambda^2 + m^2)^{\frac{1}{2}})^2}{m^2} \right) \right]$$

For $\Lambda = M_P \gg m$, **Red** is dominant :

$$\rho^{vac} \simeq \frac{\Lambda^4}{16\pi^2} \quad ; \quad p^{vac} \simeq \frac{1}{3} \frac{\Lambda^4}{16\pi^2}$$

Then : $p^{vac} \simeq \frac{1}{3} \rho^{vac}$

If matter is relativistic : $p^{mat} \simeq \frac{1}{3} \rho^{mat}$

Therefore : $p = p^{vac} + p^{mat} \simeq \frac{\rho^{vac} + \rho^{mat}}{3} = \frac{\rho}{3}$

Consequences for the CC problem

- $\rho = \rho^{vac} + \rho^{mat}$ follows evolution of relat. matter :

$$\rho(t) \propto a(t)^{-4} \quad \Leftrightarrow \quad \dot{\rho} + 3 \left(\frac{\dot{a}}{a}\right) (\rho + p) = 0$$

- Clearly this Equation also holds for ρ^{vac} and ρ^{mat} separately

$$\rho^{mat}(t) \propto a(t)^{-4} \quad \Leftrightarrow \quad \dot{\rho}^{mat} + 3 \left(\frac{\dot{a}}{a}\right) (\rho^{mat} + p_{mat}) = 0$$

$$\rho^{vac}(t) \propto a(t)^{-4} \quad \Leftrightarrow \quad \dot{\rho}^{vac} + 3 \left(\frac{\dot{a}}{a}\right) (\rho^{vac} + p^{vac}) = 0$$

(In fact, when no matter is present, $\rho = \rho^{vac} \Rightarrow$ the continuity equation holds for ρ^{vac} alone. As the ZPEs do not interact with matter, ρ^{mat} satisfies the same equation.)

- For relativistic matter :

$$\rho^{mat}(t) = \frac{\pi^2}{30} T^4 \propto a^{-4}$$

- But ρ^{vac} has the same scaling :

$$\Rightarrow \rho^{vac}(t) = \frac{\rho^{vac}(t_P)}{\rho^{mat}(t_P)} \rho^{mat}(t) \quad (t_P = \text{Planck time})$$

- From $\dot{a}/a = -\dot{T}/T \Rightarrow$ Fried. Eq., $\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3} \rho$, becomes :

$$\left(\frac{\dot{T}}{T}\right)^2 = \frac{8\pi G}{3} \left(1 + \frac{\rho^{vac}(t_P)}{\rho^{mat}(t_P)}\right) \rho^{mat}(t) = \frac{4\pi^3 G}{45} \left(1 + \frac{\rho^{vac}(t_P)}{\rho^{mat}(t_P)}\right) T^4$$

- By integrating : $T = \left(\frac{45}{16\pi^3 K G}\right)^{\frac{1}{4}} t^{-\frac{1}{2}}$

where $K = 1 + \rho^{vac}(t_P)/\rho^{mat}(t_P)$. In the standard approach:

$\rho^{vac}(t)$ is absent $\Rightarrow K = 1$

OUR MODEL

- Assume that ϕ is a sort of “**primordial field**” (out of which the other fields are born during the evolution), which describes physics at the scale $\Lambda = M_P$ at the Planck time t_P (beyond M_P ...some **UV** completion).

- Compare $\rho^{vac}(t_P) = \frac{M_P^4}{16\pi^2}$ with $\rho^\gamma(t_P) = 2 \frac{\pi^2}{30} T_P^4$ ($G = M_P^{-2} = t_P^2$) :

$$\frac{\rho^{vac}(t_P)}{\rho^\gamma(t_P)} \simeq 0.135$$

- As $\rho^{vac}(t)$ and $\rho^\gamma(t)$ have the same scaling :

$$\frac{\rho^{vac}(t)}{\rho^\gamma(t)} \simeq 0.135 \quad \text{at any time}$$

We conclude that : as **today** the ρ^γ contribution to ρ is **negligible**, the same holds true for ρ^{vac} !

- A first outcome of our model : today these red-shifting ZPEs are negligible.
- Another one : the ZPEs do not contribute to the CC as $w = 1/3$ (and not $w = -1$). They red-shift as radiation.
- This model provides a mechanism for the washing out of the ZPEs. But it has to pass an important test.
- The model predicts that ρ^{vac} is negligible today. But... what happens at the BBN epoch? Could the presence of this new radiation component screw up well tested predictions of the BBN theory?
- Let's "go back" to the BBN epoch. The total amount of radiation during this epoch is typically casted in the form :

$$\rho_R = \left(1 + \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} N_{\text{eff}} \right) \rho^\gamma$$

where (neglecting the tiny out of equilibrium neutrino contributions) $N_{\text{eff}} = 3 + \delta n$ and δn accounts for extra d.o.f. (neutrino equivalent d.o.f.), i.e. :

$$\rho_R = \left(1 + \frac{7}{8} \left(\frac{4}{11}\right)^{4/3} (3 + \delta n)\right) \rho^\gamma$$

- Comparing the **BBN** predictions with observations :

$$\delta n \leq 0.6$$

⇒ Possible additional radiation content is bounded by :

$$\rho_{\text{add}} \leq \frac{7}{8} \left(\frac{4}{11}\right)^{4/3} 0.6 \rho^\gamma = 0.136 \rho^\gamma$$

- Now compare with our previous finding

$$\rho^{\text{vac}} \simeq 0.135 \rho^\gamma \quad \text{at each time}$$

$$\text{(with } \rho^{\text{vac}}(t_P) = \frac{M_P^4}{16\pi^2} \quad ; \quad \rho^\gamma(t_P) = 2 \frac{\pi^2}{30} T_P^4 \text{)}$$

What do we learn from these results?

1. Our Model does not conflict with BBN predictions ;
2. The experimental bound ($\rho_{add} \leq 0.136 \rho^\gamma$) is saturated by our model result ($\rho^{vac} \simeq 0.135 \rho^\gamma$).

According to these findings, we might even try the following

CONJECTURE :

There is (only) one additional contribution to the radiation content coming from the ZPE of this primordial field

Brief Summary up to now :

- **Cosmic evolution** provides the mechanism to **dilute** (~ 0) the **Zero Point Energy** contribution to ρ ;
- **Zero Point Energy** **does not contribute** to the **CC** ;
- At t_P , physics is described by **one primordial quantum field**.
- **Lower Energy Theories** : **born during cosmic evolution**.

Moreover note that :

- **Low Energy Fields** : **just a convenient way to parametrize physics at lower scales (NO NEW DOF)**. \implies

⇒ When computing Vacuum Energy : do not include Zero Point Energies of Lower Energy Effective Theories. Otherwise : multiple counting of dof. Zero Point Energies of the original High Energy Theory give the whole contribution to Vacuum Energy.

ANOTHER IMPORTANT LESSON

Some Low Energy Theories have Condensates. They enter $T_{\mu\nu}$ as : $T_{\mu\nu} = \rho^{cond} g_{\mu\nu} \Rightarrow$ should give large contributions to CC.

..... BUT

Effective Field Theory scenario : no such terms. Taking into account these terms would again result in a multiple counting of dof.

Let us elucidate these points with an example

Inspired to the analysis of top condensate models by Bardeen, Hill and Lindner

- **NJL high energy theory** , defined at the scale Λ :

$$Z = \int D\bar{\psi} D\psi \exp \left[i \int d^4 x \left(\bar{\psi}(i\gamma^\mu \partial_\mu - M)\psi + \frac{g^2}{2m_0^2} \bar{\psi}\psi\bar{\psi}\psi \right) \right]$$

- **Hubbard-Stratonovic** \Rightarrow **auxiliary scalar field ϕ** :

$$Z = \frac{1}{\mathcal{N}} \int D\bar{\psi} D\psi D\phi \exp \left[i \int d^4 x \left(\bar{\psi}(i\gamma^\mu \partial_\mu - M)\psi - \frac{m_0^2}{2} \phi^2 + g\bar{\psi}\psi\phi \right) \right]$$

- **Normaliz. factor \mathcal{N}** : ensures the equality of the two Eqs.

- **Integrating high frequency modes from Λ to μ :**

$$Z = \frac{\mathcal{Q}}{\mathcal{N}} \int D\bar{\psi}_l D\psi_l D\phi_l \exp \left[i \int d^4x \left(\bar{\psi}_l (i\gamma^\mu \partial_\mu - M - \delta M) \psi_l + g \bar{\psi}_l \psi_l \phi_l + \frac{1}{2} Z_\phi \partial^\mu \phi_l \partial_\mu \phi_l - \frac{m_0^2 + \delta m_0^2}{2} \phi_l^2 - \frac{\lambda}{24} \phi_l^4 \right) \right]$$

ϕ_l and ψ_l : fields with Fourier components from 0 to μ .

- **Normally, the factor \mathcal{Q}/\mathcal{N} is not (cannot be !) considered : no knowledge of the High Energy Theory ! No effect for evaluating Green's fcts. (scattering processes). However : If we compute the Vacuum Energy from this Effective Lagrangian, we end up with a result which differs from the one obtained from the "Fundamental" = "High Energy" NJL theory unless the factor \mathcal{Q}/\mathcal{N} is taken into account. Origin of the mismatch : erroneous counting of dof.**

- **Condensates** – The same argument applies when **additional contributions** to the vacuum energy come from the appearance of **condensates** such as, for instance, a vacuum expectation value for ϕ_t .

Let me please repeat myself on this point :

The low energy theory is “**just**” a convenient way to describe physics at a lower scale in terms of a parametrization which is fit to that scale,

No new dof are created!

Curved space-time (FRW)

I have presented the model in a simple setting : **Minkowski space-time**.

We can consider our **QFT** in an **expanding universe** : a **FRW background** (which is of interest to us).

The “**adiabatic basis**” approach allows for : **mode decomposition**, definition of a **Fock space** at each time.

For the leading “**divergences**” in p^{vac} & ρ^{vac} again we have :

$$p^{vac} = \rho^{vac}/3 \quad \& \quad \rho^{vac} = a(t)^{-4} \Lambda^4$$

which are nothing but our results.

From the Effective Action

A common believe : when $T^{\mu\nu}$ is computed from the Effective Action, $T^{\mu\nu} = \frac{\delta\Gamma}{\delta g_{\mu\nu}}$, we necessarily get $w = -1$ even for the quartic divergent part of $T^{\mu\nu}$.

This gives me the opportunity to come back to a central point of my talk : The quartic divergences in $T^{\mu\nu}$, when a physical cut-off is used, come from summing up the ZPE up to the scale M_P .

In any other, formal, regularization, this very physical meaning of the quartic divergent term is lost.

Now remember : in the usual computation of Γ we use Schwinger-De Witt.

Consider the $1l$ contribution to the **Effective Action** :

$$\Gamma = -\frac{i}{2} \text{Tr} \ln[-G_F]$$

When we use the **proper time**, we invert the $\int ds$ with the $\int d^4k$ integration (we perform the $\int d^4k$ first) (*).

Don't do that ... you'll get, **of course**, the physical result.

(*)

$$\int \frac{d^4k}{(2\pi)^4} \frac{1}{k^2 - m^2 + i\epsilon} = -i \int \frac{d^4k}{(2\pi)^4} \int_0^\infty ds e^{is(k^2 - m^2 + i\epsilon)}$$

CONCLUSIONS

1. General results

- At t_P physics described by Effective Field Theory (-ies) with UV cutoff $M_P \Rightarrow$ Vacuum Energy Density ρ^{vac} undergoes a cosmic scaling : ρ^{vac} is negligible at present time t_0
- How does it come? For an Effective Field Theory : $T_{\mu\nu} = \langle\langle \hat{T}_{\mu\nu} \rangle\rangle$ is such that $p^{vac} \sim \rho^{vac}/3$
- ρ^{vac} does not contribute to CC ($w^{vac} \sim 1/3$ while $w^{CC} \sim -1$)
- Condensates do not contribute to CC (otherwise : multiple counting of DOF)

(...trying to be even more brave...if possible...)

2. Model results

- Assuming a ϕ “primordial field” at $\Lambda = M_P$ and t_P and comparing ρ^{vac} with ρ^γ : $\rho^{vac} \simeq 0.135 \rho^\gamma$
- Comparing with the **BBN** experimental bound on possible additional radiation content (neutrino equivalent d.o.f.)

$$\rho_R = \left(1 + \frac{7}{8} \left(\frac{4}{11}\right)^{4/3} (3 + \delta n)\right) \rho^\gamma \quad ; \quad \delta n \leq 0.6 :$$

$$\rho_{add} \leq \frac{7}{8} \left(\frac{4}{11}\right)^{4/3} 0.6 \rho^\gamma = 0.136 \rho^\gamma \quad \Rightarrow$$

- Our Model does not conflict with **BBN** predictions ;
- The experimental bound is saturated by our model ;
- Accordingly we have conjectured : There is (only) one additional contribution to ρ_R coming from the **ZPE** of ϕ .