

# Leptonic color models from AdS/CFT

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# Outline

- Motivation for leptonic color models
- Trinification and Quartification
- Obtaining product group GUTs from AdS/CFT
- Quartification from AdS/CFT
- Flipped Quartification

K.S. Babu, T.W. Kephart, H. Päs: [arXiv: 0709.0765](https://arxiv.org/abs/0709.0765)

# Motivation for leptonic color models

Why are quarks different from leptons:

- exist in three colors and transform as triplets under  $SU(3)_c$
- carry fractional electric charges

Assume that this difference is a low energy artefact:

- Assume leptonic color  $SU(3)_l$  symmetry
- Breaking  $SU(3)_l \rightarrow SU(2)_l$  gives unconfined lepton states and explains electric charge differences between leptons and quarks

Foot, Lew 1990; Foot, Lew, Volkas 1991

→ Possibility to unify  $SU(3)_l$  and SM gauge group into the quartification gauge group

$$SU(3)_L \times SU(3)_R \times SU(3)_c \times SU(3)_l$$

Joshi, Volkas, 1992; Babu, Ma, Willenbrock, 2003

# Trinification models

Quartification models are a generalization of Trinification models:

$$SU(3)_c \times SU(2)_L \times U(1)_Y \rightarrow SU(3)_c \times SU(3)_L \times SU(3)_R$$

De Rujula, Glashow, Georgi, 1984

→ Aesthetic: fermions in (bi-)fundamental representations:

$$3[(3\bar{3}1) + (13\bar{3}) + (\bar{3}13)]$$

with

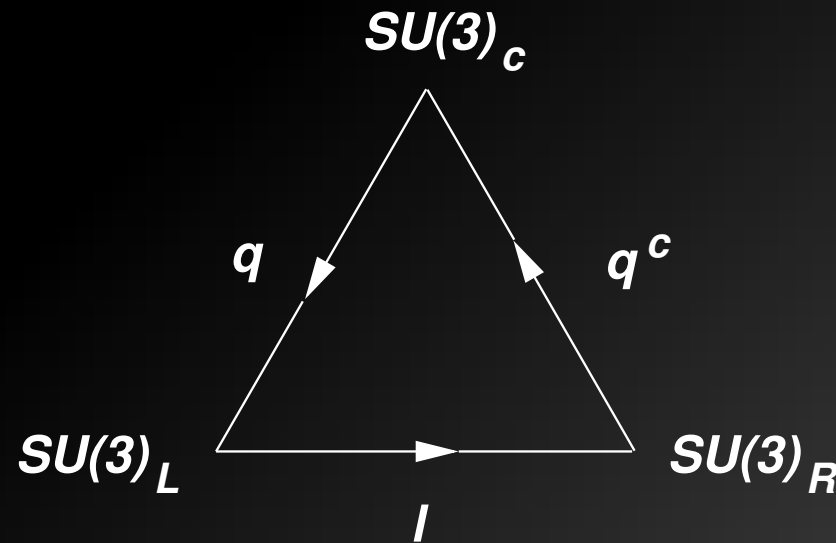
$$l \sim (1, 3, \bar{3}), \quad q \sim (3, \bar{3}, 1), \quad q^c \sim (\bar{3}, 1, 3)$$

→ Attractive: suppressed proton decay rate (only due to Higgs exchange)

Babu, He, Pakvasa, 1986

# Fermion content of Trinification models

The fermion content can be arranged into a **quiver diagram**:



In terms of  $SU(3)_L \times SU(3)_R$  one gets:

$$l \sim \begin{pmatrix} N & E^c & \nu \\ E & N^c & e \\ \nu^c & e^c & S, \end{pmatrix}$$

rows:  $I_{3L} = (1/2, -1/2, 0)$  and  $Y_L = (1/3, 1/3, -2/3)$

columns:  $I_{3R} = (-1/2, 1/2, 0)$  and  $Y_R = (-1/3, -1/3, 2/3)$

# Fermion content of Trinification models

$$q \sim \begin{pmatrix} d & u & h \\ d & u & h \\ d & u & h \end{pmatrix}, \quad q^c \sim \begin{pmatrix} d^c & d^c & d^c \\ u^c & u^c & u^c \\ h^c & h^c & h^c \end{pmatrix}$$

columns in  $q$ :  $I_{3L} = (-1/2, 1/2, 0)$ ,  $Y_L = (-1/3, -1/3, 2/3)$

rows in  $q^c$ :  $I_{3R} = (1/2, -1/2, 0)$ ,  $Y_L = (1/3, 1/3, -2/3)$

Electric charge operator:

$$Q = I_{3L} - \frac{Y_L}{2} + I_{3R} - \frac{Y_R}{2} = I_{3L} + \frac{Y}{2}$$

- Exotic fermions  $h(h^c)$ ,  $E(E^c)$ , and  $N, N^c, S$ : charges  $\mp 1/3, \mp 1, 0$
- Trinification models can be embedded into  $E_6$  but have no quark-lepton-symmetry!

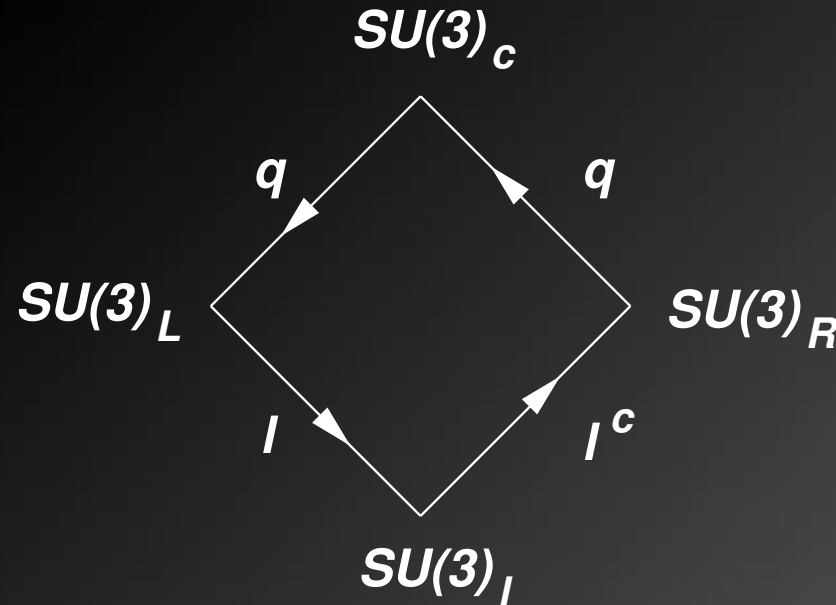
# Quartification models

Gauge group:

$$SU(3)_l \times SU(3)_L \times SU(3)_c \times SU(3)_R$$

with fermions in bi-fundamental representations:

$$3(\bar{3}\bar{3}11) + (13\bar{3}1) + (113\bar{3}) + (\bar{3}113)$$



and electric charge operator:

$$Q = I_{3L} - \frac{Y_L}{2} + I_{3R} - \frac{Y_R}{2} - \frac{Y_l}{2}$$

# Fermion content of Quartification models

Lepton content:

$$l \sim \begin{pmatrix} x_1 & x_2 & \nu \\ y_1 & y_2 & e \\ z_1 & z_2 & N \end{pmatrix}, \quad l^c \sim \begin{pmatrix} x_1^c & y_1^c & z_1^c \\ x_2^c & y_2^c & z_2^c \\ \nu^c & e^c & N^c \end{pmatrix}$$

$\underline{l}$  columns:  $Y_l = (-1/3, -1/3, 2/3)$  rows:  $I_{3L} = (1/2, -1/2, 0)$ ,  $Y_L = (1/3, 1/3, -2/3)$

$\underline{l}^c$  rows:  $Y_l = (1/3, 1/3, -2/3)$ , columns:  $I_{3R} = (-1/2, 1/2, 0)$ ,

$Y_L = (-1/3, -1/3, 2/3)$

## New particles:

- $N$  and  $N^c$ : neutral
- $SU(2)_l$  doublet leptons  $(x, y, z)$ : charges  $(1/2, -1/2, 1/2)$   
 $(x^c, y^c, z^c)$ : charges  $(-1/2, 1/2, -1/2)$

Interesting feature: gauge coupling unification without supersymmetry

Babu, Ma, Willenbrock, 2003

# Product group GUTs from AdS/CFT

Maldacena 1997: IIB superstring theory compactified on  $AdS \times S^5$  is dual to N=4 SUSY CFT

Frampton 1998:  $\mathcal{N} = 4$  SUSY can be broken down to  $\mathcal{N} = 0, 1$  SUSY by orbifolding  $S^5 \rightarrow S^5/\Gamma$  where  $\Gamma$  is a discrete group:

$\Gamma \subset SU(4) \sim O(6)$  of  $S^5 \Rightarrow \mathcal{N} = 0$

$\Gamma \subset SU(3) \subset O(6)$  of  $S^5 \Rightarrow \mathcal{N} = 1$

→ remnants of conformality property can help to solve the hierarchy problem

Kachru, Silverstein 1998; Frampton, Vafa, Kephart, Mohapatra, ...:  
AdS/CFT/Conformality model building

Choose  $N = 3$  in large  $N$  expansion and orbifold with  $Z_n \Rightarrow$  gauge group  $SU(3)^n$

Fermions: 4 of  $SU(4)$  isometry must be complex for chiral fermions

Scalars:  $6 = (4 \otimes 4)_{\text{antisym}}$  of  $SU(4)$  isometry must be real for consistent embedding  
(Frampton, Kephart, 2004)

Kephart, Päs, 2001; Kephart, Päs, 2004:

→ systematic scans of  $\mathcal{N} = 0$  and  $\mathcal{N} = 1$  SUSY models from  $\Gamma = Z_n$  orbifolds  
check for chiral fermion content and consistent with phenomenology

# Leptonic color models from AdS/CFT

To obtain

$$SU(3)_l \times SU(3)_L \times SU(3)_c \times SU(3)_R$$

with all four coupling constants equal:

→ Look for  $Z_n$  orbifold with  $n = 4z$ ,  $z \in \mathbb{Z}$

⇒  $SU(3)^4, SU(3)^8, \dots, SU(3)^{4z} \rightarrow SU(4)^n$

with each factor being the diagonal subgroup of  $z$  of the original groups

n=4: no viable model with chiral fermions

n=8: 24 possible symmetry breaking patterns - two possible models:

- one full quartifications model:  $4 = (\alpha^1, \alpha^1, \alpha^3, \alpha^3)$

$$3[(\bar{3}\bar{3}11) + (13\bar{3}1) + (113\bar{3}) + (\bar{3}113)]$$

- one hybrid trinification-quartification model:  $4 = (\alpha^1, \alpha^4, \alpha^5, \alpha^6)$

$$2[(\bar{3}\bar{3}1) + (13\bar{3}) + (\bar{3}13)] + (\bar{3}\bar{3}11) + (13\bar{3}1) + (113\bar{3}) + (\bar{3}113)$$

# Quartification from Ads/CFT

Fermions:  $4 = (\alpha^1, \alpha^1, \alpha^3, \alpha^3) \rightarrow$

$$2[(3\bar{3}11111) + (311\bar{3}111)]_F + \text{cyclic permutations}$$

Scalars:  $6 = (4 \otimes 4)_{\text{antisym}}$

$$[(31\bar{3}1111) + 4(3111\bar{3}11) + (311111\bar{3}1) + h.c.]_S + \text{cyclic permutations}$$

$SU(3)^8 \rightarrow SU(3)^4$  by assigning VEVs to

$$(31\bar{3}1111)_S, (131\bar{3}111)_S, (111131\bar{3}1)_S, (1111131\bar{3})_S$$

$\Rightarrow$  **chiral  $SU(3)^4$  Fermions:**

$$4[(3\bar{3}11) + (13\bar{3}1) + (113\bar{3}) + (\bar{3}113)]_F$$

and **scalars**

$$10[(31\bar{3}1) + (131\bar{3}) + h.c.]_S + 2[(8111) + (1811) + (1181) + (1118)]_S$$

**problem:** few scalars, SSB to the SM requires leptonic color condensate, difficult for small  $\Lambda_{ICD}$

# Minimal leptonic color from AdS/CFT

Fermions:  $4 = (\alpha^1, \alpha^4, \alpha^5, \alpha^6) \rightarrow$

$$(3\bar{3}111111)_F + (3111\bar{3}111)_F + (31111\bar{3}11)_F + (311111\bar{3}1)_F + \text{cyclic permutations}$$

Scalars:  $6 = (4 \otimes 4)_{\text{antisym}}$

$$(3\bar{3}111111)_S + (31\bar{3}11111)_S + (311\bar{3}1111)_S + (31111\bar{3}11)_S + (311111\bar{3}1)_S \\ + (3111111\bar{3})_S + \text{cyclic permutations}$$

**SSB**  $SU(3)^8 \rightarrow SU(3)^4$  by assigning VEVs to

$$(3\bar{3}111111)_S, (11311\bar{3}11)_S, (111311\bar{3}1)_S, (1111311\bar{3})_S$$

$\Rightarrow$  chiral  $SU(3)^4$  Fermions:

$$2[(13\bar{3}1) + (113\bar{3}) + (1\bar{3}13)]_F + [(3\bar{3}11) + (13\bar{3}1) + (113\bar{3}) + (\bar{3}113)]_F$$

# Minimal leptonic color from AdS/CFT

⇒ Scalars:

$$3[(3\bar{3}11) + (13\bar{3}1) + (113\bar{3}) + (\bar{3}113) + h.c.]_S \\ + 4[(31\bar{3}1) + (131\bar{3}) + h.c.]_S + 2[(8111) + (1811) + (1181) + (1118)]_S$$

- VEV to a  $(3\bar{3}11)$ : 3 family trinification model
- VEVs are given to two  $(131\bar{3})$  representations: **realistic quark masses and non-zero mixing angles**, VEVs to  $(3\bar{3}11)$  and to  $(311\bar{3})$ :  $SU(3)_l \rightarrow SU(2)_l$
- **do we still have gauge coupling unification?**

# Gauge coupling unification

RGEs:

$$\frac{1}{\alpha_i(\mu)} - \frac{1}{\alpha_i(\mu')} = \frac{b_i}{2\pi} \ln\left(\frac{\mu'}{\mu}\right)$$

where  $b_n$  are the one-loop beta-function coefficients:

$$b_3 = -11 + \frac{4}{3}N_g,$$
$$b_2 = -\frac{22}{3} + 2N_q + \frac{4}{3}N_t + \frac{1}{6}N_H$$
$$b_1 = \frac{13}{9}N_q + \frac{4}{3}N_t + \frac{1}{12}N_H$$

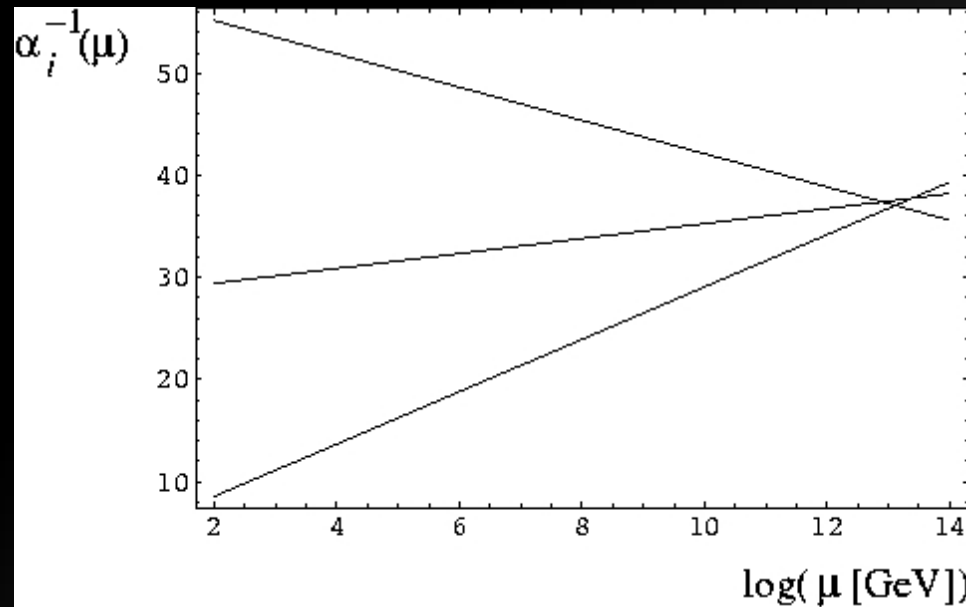
$N_g = 3$ : number of generations

$N_q = 1$ : number of quartification families

$N_t = 2$ : number of trinification families

$$\sin^2 \theta_W(M_{GUT}) = \sum I_{3L}^2 / \sum Q^2 = 9/16$$

# Gauge coupling unification



12 standard Higgs doublets of  $SU(2)_L$ : 4 Higgs doublets  $\rightarrow$  unification at  $10^{13}$  GeV  
 $\rightarrow$  chance to discover multiple Higgs doublets at the LHC

Coupling of the unbroken  $SU(2)_l$ :

$$b_{2l} = -\frac{22}{3} + \frac{4}{3}N_q.$$

$\Rightarrow \alpha_{2l}^{-1}(M_Z) \simeq 13$  between  $\alpha_{\text{weak}}$  and  $\alpha_s$ , similar phenomenology as in Babu, Ma, Willenbrock, 2003 albeit with higher  $\Lambda_{LCD} \sim MeV$

# Constraints and consistency

- Hemion masses: quartification scale hemion masses: forbidden by the  $Z_8$  orbifold symmetry  
TeV scale hemion masses: generated by adding non-renormalizable operators suppressed by  $M_{\text{Planck}}$
- Electroweak precision data: Singlet and vectorlike symmetry breaking products will not affect EW data  
hemions: vector-like under the SM gauge group, EW effects suppressed by hemion masses  
Flavor changing neutral currents?
- Proton decay: not be mediated by gauge bosons, could be induced via couplings to the extended scalar sector
- Stickballs glueballs of leptonic color SU(2): cold dark matter candidate

# Neutrino masses

SSB provides several singlet fermions:

e.g. from the  $(\bar{3}\bar{3}111111)$  representation after assigning VEV to  $(\bar{3}\bar{3}111111)$  scalars  
→ makes AdS/CFT inspired model superior to the simple quartification model, as a seesaw mechanism

$$m_\nu \sim \frac{m_{\bar{\nu}\nu}^2}{M_{\bar{3}\bar{3}111111}}$$

can be implemented to generate light neutrino masses without adding right-handed neutrinos by hand

→ interesting neutrino sector: multidimensional neutrino mass matrices, seesaw II?

→ Leptogenesis,....

# Flipped Quartification

Quartification models: possible to non-trivially flip order of factor groups

$$SU(3)_L \times SU(3)_R \times SU(3)_c \times SU(3)_l \rightarrow SU(3)_c \times SU(3)_L \times SU(3)_R \times SU(3)_l$$

SSB:

- $SU(3)_L \rightarrow SU(2)_L \times U(1)_A$
- $SU(3)_R \rightarrow U(1)_B \times U(1)_C$
- $SU(3)_l \rightarrow SU(2)_l \times U(1)_D$

$$Q = T_3 - \frac{1}{6}A + \frac{1}{2}B - \frac{1}{6}C + \frac{1}{3}D$$

yields all necessary third family states in  $SU(2)_l$  singlets except for  $(b_R)^c$

→ forms bound state with some  $(1, 1, 2)$  scalar

→ interesting consequences for  $b$ -physics

Dent, Kephart, Päs, work in progress

# Summary and Conclusions

- Models with the **quartification gauge group**  $SU(3)^4$  are attractive, since they allow for a **quark-lepton symmetry**
- **Interesting variants** of such models can be **obtained from  $Z_8$  orbifolded AdS/CFT**
- **Gauge couplings unify**
- **EW precision, proton decay, neutrino masses, dark matter ok**
- **Interesting  $b$  physics** anticipated in “flipped” variants
- **New scalars at the LHC**